Nonstoichiometric Alkali Ferrites and Aluminates in the Systems NaFeO₂-TiO₂, KFeO₂-TiO₂, KAlO₂-TiO₂, and KAlO₂-SiO₂

C. LI, A. F. REID, AND S. SAUNDERS

Division of Mineral Chemistry, CSIRO, P.O. Box 124, Port Melbourne, Victoria 3207, Australia

Received March 2, 1971

Examination of NaFeO₂-TiO₂ compositions has shown that the α -NaFeO₂ structure can be given a range of nonstoichiometry Na_{1-x}Fe_{1-x}Ti_xO₂, $O \le x \le 0.28$. The additional Ti atoms are randomized with Fe atoms, and charge balance is maintained by a partial occupancy of the Na atom layers. Fractional occupancy of alkali ion sites in other compounds in the NaFeO₂-TiO₂ system and in the systems KAIO₂-TiO₂, KAIO₂-SiO₂, and KFeO₂-TiO₂ was also found. In TiO₄ tetrahedra, Ti⁴⁺ was deduced to have a Ti-O bond length of 1.92 Å.

Introduction

The combination of alkali metal oxides with the oxides of trivalent metal ions provides a large number of ABO₂ compounds (1-3) of which α -NaFeO₂ (4), based on the rocksalt structure, and KFeO₂ (5), based on the cristobalite structure, are typical.

We have found that a number of ABO_2 compounds can incorporate TiO_2 in their structures, and additionally, those based on cristobalite can incorporate SiO₂. The details of the resulting nonstoichiometry and its relationship to that of other nonstoichiometric phases in the ABO_2 -TiO₂ systems form the subject of the present paper.

Experimental

Sodium based compositions of the type $Na_{1-x}Fe_{1-x}Ti_xO_2$ were prepared from analytical reagent grade sodium oxalate, and Fisher certified Fe_2O_3 and TiO₂. After thorough grinding, the preparations were heated at 600°C for several hours, and then at 650°C for x = 0, 800°C for x = 0.05 and 0.10, and at 1000°C for x > 0.15. The α form of NaFeO₂ is transformed near 760° to a β form (6-8) based on the wurtzite structure. In the present work, this transformation temperature was found to be raised above 1000°C by addition of between 10 and 20 mole % TiO₂ to α -NaFeO₂ and the temperatures of preparation were chosen to avoid the α - β conversion.

Potassium-based compounds were made from dried potassium carbonate and, variously, Al_2O_3 , Fe_2O_3 , TiO_2 and precipitated silica. After initial grinding and firing at 900°C, all preparations based on KAIO₂ were reground and heated overnight at 900°C. KFeO₂ was prepared at 800°C, and $K_{1-x}Fe_{1-x}Ti_xO_2$ compositions at 750°C. The potassium compounds were sensitive to atmospheric moisture, and were cooled and manipulated in a dry atmosphere.

Lithium compounds were made from lithium carbonate dried at 300°C and were heated at a final temperature of 1000°C.

X-Ray Data

X-Ray data were obtained on a powder diffractometer fitted with monochromator and proportional counter. Recordings were made at slow scan rates, and either silicon powder, $a_0 = 5.4301$ Å or potassium chloride, $a_0 = 6.293$ Å, was included as an internal standard. Potassium compounds were protected by an atmosphere of dry nitrogen as it was found that absorption of atmospheric moisture markedly affected the cubic lattice parameter. For the hexagonal compounds, final lattice parameters were obtained by a least squares fit of the observed data. Reflection intensities for Na_{0.75}Fe_{0.75}Ti_{0.25}O₂ were obtained from the areas of recorded diffraction peaks.

The lower limit of the narrow fractional occupancy

range for sodium in the NaFeTiO₄ structure, 0.970 \pm 0.005, was determined by measuring the relative intensities of the NaFeTiO₄ (9) and Na_xFe_xTi_{2-x}O₄ (10) powder patterns obtained at several compositions in their region of coexistence, and extrapolating to a zero value for Na_xFe_xTi_{2-x}O₄.

Results

The variation in lattice parameters for α -NaFeO₂-TiO₂ solid solutions are shown in Fig. 1 and limiting values are given in Table I. It is evident that the composition range Na_{1-x}Fe_{1-x}Ti_xO₂, exists with $0 \le x \le 0.28$. Beyond this concentration, the lattice parameters remained constant, and NaFeTiO₄ (9), the next phase in the pseudobinary NaFeO₂-TiO₂ system, began to appear.

Structure factor calculations for Na_{0.75}Fe_{0.75} Ti_{0.25}O₂ were made after converting the observed intensities to structure factors, with appropriate allowance for multiplicity and Lorentz polarization. Calculations were based on the hexagonal cell with a = 3.000 and c = 16.410 Å. Within the space group

FIG. 1. Lattice parameters of $Na_{1-x}Fe_{1-x}Ti_xO_2$ as a function of $NaFeO_2$: TiO₂ ratio.

TABLE I

LATTICE PARAMETERS FOR LIMITING COMPOSITIONS

Compound	Lattice par		
	<i>x</i> = 0	$x = x_{\text{limit}}$	Limit of <i>x</i>
$Na_{1-x}Fe_{1-x}Ti_xO_2$	a = 3.020(3)	a = 2.996(3)	0.28
$\mathbf{K}_{1-x}\mathbf{A}\mathbf{l}_{1-x}\mathbf{T}\mathbf{i}_{x}\mathbf{O}_{2}$	c = 16.10(1) a = 7.724(5)	c = 16.40(1) a = 7.820(5)	0.25
$K_{1-x}Fe_{1-x}Ti_xO_2$	a = 7.954(5)	a = 7.993(5)	0.20
$K_{1-x}Al_{1-x}Si_xO_2$	<i>a</i> = 7.724(5)	a = 7.702(8) c = 7.646(8)	0.20

"Numbers in parentheses represent three standard deviations, e.g. 3.020 ± 0.003 .

 $R\overline{3}m(I)$, sodium atoms with a fractional occupancy of 0.75 were placed at the origin, with (0.75Fe + 0.25Ti) randomized in point position (3b) at 00½. Oxygen was placed at 00z in (6c), with an initial value of z of 0.250. After several cycles of leastsquares refinement, a final value of z of 0.2309 ± 0.0021 was obtained, with $R_{F_0} = 5.0\%$. Attempts to allow a fractional occupancy for the Fe, Ti sites, with incorporation of Fe or Ti in the Na layers, gave greatly increased values of R_{F_0} . Observed and calculated powder intensity data, for which $R_{I_0} =$ 6.9%, are shown in Table II.

TABLE II

X-RAY POWDER DIFFRACTION PATTERN OF Na_{0.75}Fe_{0.75}Ti_{0.25}O₂^{*a*}; Cu K_{α_1} , KCI Internal Standard

h k l	$\sin^2 \theta_{obsd}$	$\sin^2 \theta_{calcd}$	$I_{\rm obsd}$	I b caled	
003	0.01980	0.01985	274	280	
006	0.07930	0.07938	50	39	
101	0.09012	0.09008	56	51	
012	0.09657	0.09669	70	65	
104	0.12316	0.12315	324	337	
107	0.19604	0.19592	54	47	
018	0.22914	0.22899	86	83	
110	0.26380	0.26361	83	92	
113	0.28356	0.28346	39	33	
1.0.10	0.30842	0.30837	23	22	
116	0.34309	0.34299	24	28	
021)	0.05465	0.35368	7	6	
0.1.11	0.35465	0.35468	17	16	
202	0.36009	0.36030	9	11	
024	0.38676	0.38676	33	38	
208	0.49258	0.49260	21	19	
1.1.12	0.58115	0.58113	19	22	

" $a = 3.000 \pm 0.005$, $c = 16.40 \pm 0.01$ Å.

^b Absolute units $\times 10^{-2}$.





FIG. 2. Lattice parameters of $K_{1-x}AI_{1-x}Ti_xO_2$, $K_{1-x}AI_{1-x}Si_xO_2$, and $K_{1-x}Fe_{1-x}Ti_xO_2$ as a function of x. For the last compound, values beyond x = 0.9 are the cubic lattice parameter corresponding to the unit cell volume of the tetragonal cell.

In the KAIO₂-TiO₂ and KFeO₂-TiO₂ systems, Fig. 2, the lattice parameter changes indicated that solid solution $K_{1-x}Al_{1-x}Ti_xO_2$ with 0 < x < 0.25and $K_{1-x}Fe_{1-x}Ti_xO_2$ with 0 < x < 0.20 can be formed. For each system, the limiting value of x was also inferred from the composition at which a second phase began to appear. The $K_{1-x}Fe_{1-x}Ti_xO_2$ compositions were unstable above temperatures of 800–900°C, depending on Ti content. Single-phase preparations were obtained at temperatures sufficiently low to prevent decomposition, and typically 750°C.

In the KAlO₂-SiO₂ system, SiO₂ entered the cristobalite-based KAlO₂ lattice up to the composition K_{0.80}Al_{0.80}Si_{0.20}O₂; beyond this composition the tridymite-based KAlSiO₄ structure (11-13) began to appear. The pseudocubic lattice parameters for K_{1-x}Al_{1-x}Si_xO₂ are included in Fig. 2; beyond x = 0.90 a pronounced splitting of the cubic lines, particularly (220), was evident, indicative of a tetragonal distortion. The powder patterns did not contain enough reflections for unambiguous indexing, but were consistent with cell parameters

 $a_0 = 7.702$, $c_0 = 7.646$ Å at a limiting composition x = 0.20.

LiCrO₂ (14), although of the α -NaFeO₂ type, showed little or no tolerance for accommodation of TiO₂, and no lattice parameter change could be detected when LiCrO₂-TiO₂ starting mixtures were reacted. Instead, the spinel LiCrTiO₄ (15) appeared, and the amount of LiCrO₂ progressively decreased.

Discussion

NaFeO₂-TiO₂ System

The α -NaFeO₂ structure (1, 4) is derived from that of rocksalt, with an ordering of Na and Fe into layers normal to the original cube diagonal. The resulting rhombohedral distortion produces a structure which contains layers of Fe–O₆ octahedra separated by layers of Na atoms (16). In the hexagonal setting, a = 3.020 Å and c = 16.10 Å, Table I, all atoms lie in planes normal to the c axis. Solid solution of α -NaFeO₂ type compounds with others also based on the rocksalt-type lattice have been observed, e.g., between LiCrO₂ and MgO (17) and between KPrO₂ and K₂PrO₃ (18), but in such cases structural sites are not left vacant.

For $Na_{0.75}Fe_{0.75}Ti_{0.25}O_2$, the structure factor calculations showed that it is the sodium layers which contain the vacancies. As might be expected, the transition metal octahedral layers remain fully occupied. The lattice parameter plots, Fig. 1, show that the structure contracts in the *a* axis direction and expands along the c axis (thus increasing the rhombohedral distortion) as the sodium occupancy falls. The metal-oxygen distances for several α -NaFeO₂ isotypes having similar unit cell dimensions but containing metal ions of significantly different sizes (19) are shown in Table III. These distances show a very narrow range, indicating that in the metal-oxygen octahedral layers, the M-O separations are constrained to a relatively fixed value by the packing of the oxygen atoms. The Na-O distances show a slightly greater variability, and between α -NaFeO₂ and Na_{0.75}Fe_{0.75}Ti_{0.25}O₂ there appears to be a small but real increase in Na-O separation in addition to the increased rhombohedral distortion of the Na-O₆ octahedron.

Partial occupancy of the alkali ion layers is presumably common for the α -NaFeO₂ type, provided that the ions which have to randomize in the octahedral layer sites are roughly comparable in size and the alkali ion is not Li⁺, which can itself randomize with Ti⁴⁺, Cr³⁺, or Fe³⁺. In the system NaTiO₂-TiO₂, e.g., there should exist a composition range Na_xTiO₂-TiO₂ with $1 \ge x > 0.70$ based on

TABLE	111

STRUCTURAL PARAMETERS FOR SOME α -NaFeO₂ Isotypes

	<i>a</i> (Å) ± 0.003	c (Å) ± 0.01	V (Å) ³	Ζ	Na–O (Å)	MO (Å)
Na _{0.75} Fe _{0.75} Ti _{0.25} O ₂	3.000	16.40	127.8	0.2309(21)	2.413(18)	2.027(16)
NaFeO ₂	3.020	16.10	127.2	0.2309ª	2.3994	2.027*
Na ₂ CoTiO ₄ ^b	3.008	16.18	126.8	0.2314(60)	2.394(56)	2.030(56)
Na ₂ NiTiO ₄ ^b	3.010	16.09	126.2	0.2320(30)	2.379(29)	2.029(29)

^a Assuming the same fractional atomic coordinate as for Na_{0.75}Fe_{0.75}Ti_{0.25}O₂.

* Ref. (20).

the NaTiO₂ (21) end member. It is also possible that sulphides and selenides with the α -NaFeO₂ structure can exhibit a similar fractional occupancy of the alkali ion sites when charge compensated with, for example, the disulphides of Ti, Nb, or Mo. Indeed, a wide range of occupancy for potassium in K_xTiS₂ has been observed, with the α -NaFeO₂ structure being retained for $1 \ge x > 0.28$ (22).

The α - β -NaFeO₂ transition temperature was considerably raised by incorporation of TiO₂ in the structure, being greater than 1000°C for Na_{0.8}Fe_{0.8}Ti_{0.2}O₂ as compared with 760°C for NaFeO₂ (7). This effect can be ascribed to the preference of Ti⁴⁺ for octahedral sites, as observed, e.g., in garnets (23-25). While it is possible that Ti⁴⁺ could enter the tetrahedral sites in the wurtzite-like form of β -NaFeO₂, at a nominal composition Na_{0.9}Fe_{0.9}Ti_{0.1}O₂ both α and β forms coexisted at 1000°C, and the lattice parameters of the α phase showed it to be Ti-rich with respect to the β phase.

In the pseudobinary NaFeO₂-TiO₂ system there are three further sodium-based phases, each of which can exhibit a range of partial occupancy of sodium ions. In each case, the crystallographic sodium sites are formed within a fully occupied framework of metal-oxygen octahedra containing randomized Fe^{3+} and Ti^{4+} . In NaFeTiO₄ (9), which is of calcium ferrite type (26–28), Na atoms occupy sites 3 Å apart in infinite linear tunnels, with a lower occupancy limit of 0.970 \pm 0.005. In Na_xFe_xTi_{2-x}O₄, $0.90 \ge x \ge 0.75$ (10), Na atoms partially occupy double tunnels of staggered sites 3 Å apart and in the NaFeTi₃O₈ structure (29-31) isotypic with $Na_{x}Ti_{4}O_{8}$ (32), the sodium atom occupancy can decrease to 0.67, giving a composition Na_{0.67}Fe_{0.67} $Ti_{3.33}O_8$ (30). In this structure sodium ions are contained in linear tunnels of perovskite-like sites 3.8 Å apart.

The NaFeO₂-TiO₂ tieline contains the only ternary oxide compounds so far known in the

Na₂O-Fe₂O₃-TiO₂ system, shown in Fig. 3. Attempts to prepare other ternary oxides have so far produced only mixtures of those shown. In this system, the binary oxide compounds $Fe_2Ti_3O_9$



FIG. 3. The ternary oxide system Na₂O-Fe₂O₃-TiO₂. Of the compounds on the Na₂O-TiO₂ tieline only Na₂Ti₃O₇ (36) and $Na_2Ti_6O_{13}$ (37) have been structurally characterized. Na₂TiO₃, melting point 1030°C, has been studied as an end member of a number of phase diagrams (38), and is believed to be similar in structure to Na_2SnO_3 (39). The structure of $Na_8Ti_5O_{14}$ (38, 40) is unknown. A compound $Na_4Ti_3O_8$ is also believed to have a transient existence (41). On the Fe₂O₃-TiO₂ tieline, Fe₂Ti₃O₉ (34) is similar to but not identical with Cr₂Ti₂O₇ (42). A naturally occurring form of Fe₂Ti₃O₉ (32) based on the ilmenite structure has also been reported with the Fe²⁺ cation layers of ilmenite replaced by layers in which Fe³⁺ ions occupy two-thirds of the original Fe²⁺ sites. The structure of Fe₂TiO₅, pseudobrookite, has long been known (43). On the Fe_2O_3 -Na₂O tieline, Na₃Fe₅O₉ has been described by Romers et al. (35). The compounds on the NaFeO₂-TiO₂ tieline are discussed in the text.

(33, 34) and Na₃Fe₅O₉ (35), have themselves been characterized only recently.

KFeO₂ and KAlO₂ with TiO₂

The structures of $KAlO_2$ and $KFeO_2$ (2, 44) are based on that of β -cristobalite, the highest temperature, lowest density modification of SiO_2 (13, 45). In the idealized form of this cubic structure, Si atoms occupy corners, face centers and the four $\frac{1}{4}$ $\frac{1}{4}$ positions of the unit cube of side 7.05 Å. Each Si is at the center of four neighbors as in the arrangement of carbon atoms in diamond, with the addition of oxygen atoms midway between silicon atoms. The real structure (13, 44-46) is a distortion of this, involving the less symmetrical space group T^4 $(P2_13)$. The center of the unit cube contains a large unoccupied cage site, as shown in Fig. 66 of Ref. (13), and in KAlO₂ and KFeO₂ this is occupied by potassium ions, with Fe or Al ions occupying the Si sites. This arrangement gives K atoms a 12-fold coordination, with bonds to two oxygen atoms of each of the Al-O₄ or Fe-O₄ tetrahedra situated at the cube faces. The oxygen coordination number is correspondingly increased from two to three. As for β -cristobalite, the real structure must involve a distortion of atoms from these ideal positions, otherwise the Al-O, Fe-O, and K-O distances would be unrealistic (19, 45). However, detailed atomic coordinates have not yet been determined.

For KAlO₂ and KFeO₂, nonstoichiometry produced by inclusion of TiO₂ in the lattice extends to the compositions $K_{1-x}Al_{1-x}Ti_xO_2$ with $0 \le x \le 0.25$, and $K_{1-x}Fe_{1-x}Ti_xO_2$ with 0 < x < 0.20. On the basis of the large size of potassium atoms and their position in cage-like sites in a cristobalite framework, this nonstoichiometry will certainly involve a partial occupancy of the cage sites. The required charge balance is then retained by a randomization of AI^{3+} and Ti⁴⁺ or Fe³⁺ and Ti⁴⁺ in the framework of metal-oxygen tetrahedra. Tetrahedral coordination to oxygen is unusual for Ti⁴⁺, although 5-coordination in the form of square pyramids in $Ba_2 TiSi_2O_8$, fresnoite (47), and trigonal bipyramids in $K_2Ti_2O_5$ (48) have been reported. However, Espinoza (23) has shown that in particular garnets, Ti⁴⁺ can replace Ga^{3+} in up to one sixth of the tetrahedral sites, and lesser amounts of Ti4+ in tetrahedral sites have been observed in other garnets (24).

Substitution of Ti^{4+} for Fe^{3+} in the tetrahedral sites of KFeO₂ *increases* the lattice parameter, as shown in Fig. 2, and it is of interest to try to estimate the size of the Ti-O₄ tetrahedron. For threecoordinate oxygen atoms, Si-O, Al-O, and Fe-O tetrahedral bond lengths are 1.62, 1.75, and 1.85 Å,

respectively, (19, 49, 50) and a plot of unit cell volume vs average tetrahedral bond length (51) for the three compositions $K_{0.8}Al_{0.8}Si_{0.2}O_2$, $KAlO_2$, and KFeO₂ was found to be linear. Using this line as a basis, the change in lattice parameter as Ti⁴⁺ replaces Al^{3+} in $K_{1-x}Al_{1-x}Ti_xO_2$ (Fig. 2) corresponds to a tetrahedral Ti-O bond length of 1.92 Å. This is abnormally high in relation to typical Ti-O octahedral distances of 2.00 Å (19) and indicates that a tetrahedral Ti-O₄ configuration is not only unstable, but also relatively large. Substantial concentrations can be expected to exist only in a relatively unconstrained structure such as that of KFeO₂ where the framework can expand to match the tetrahedral ion size, and randomization of Ti⁴⁺ with Si⁴⁺ in tetrahedral silicate sites is unlikely. Thus in the titanium-substituted andradite garnets $Ca_3Fe_2Si_3O_{12}$ (52), e.g., in which up to 50% of the Si⁴⁺ can be replaced by Ti⁴⁺, Mössbauer spectroscopy shows that an occupancy rearrangement occurs with Ti⁴⁺ entering octahedral sites and Fe³⁺ randomizing with Si⁴⁺ in the tetrahedral sites.

When further TiO_2 is added to the KAlO₂ and KFeO₂-based phases, a second phase (53) typified by $K_{0.75}Fe_{0.75}Ti_{1.25}O_4$ appears, and disproportionation to give this phase also occurs for $K_{1-x}Fe_{1-x}Ti_{x}O_{4}$ compositions if they are heated to too high a temperature. K_{0.75}Fe_{0.75}Ti_{1.25}O₄ contains puckered layers of (Fe, Ti)-O₆ octahedra, with K atoms partially occupying the sites in a structure which is related to that of $Rb_rMn_rTi_{2-r}O_4$ (54). As the Ti:Fe or Ti:Al ratio is further increased, the hollandite structure, in which potassium atoms are contained in linear tunnels bounded by a framework of edge-shared pairs of octahedra, (55) is formed. The hollandite isotype of end member composition KAlTi₃O₈ has an existence range down to $K_{0.75}Al_{0.75}Ti_{3.25}O_8$ (56), and we have observed a similar occupancy in potassium iron titanate of the hollandite form.

KAlO₂-SiO₂

As expected for a structure based on cristobalite, KAlO₂ can accommodate SiO₂ in the framework sites up to a composition $K_{1-x}Al_{1-x}Si_xO_2$ with $0 \le x \le$ 0.20. Presumably since the Si:Al ratio is relatively low, and the lattice parameter change is linear with Si concentration, Al and Si are randomized in these sites. The decrease in unit cell volume as Si replaces Al, Fig. 2, is consistent with the unit cell volume vs tetrahedral bond length values observed for KAlO₂ and KFeO₂, assuming an Si-O bond length of 1.62 Å (19, 49).

Tridymite, the next most dense silica polymorph,

also contains large empty sites, but only half the number of those in cristobalite. The sites lie in linear tunnels formed by rings of six tetrahedra corner joined with alternate tetrahedra pointing up and down. The compound KAlSiO₄, kalsilite (12), is based on this tridymite framework and in the KAlO₂-SiO₂ system, KAlSiO₄ was found to co-exist with $K_{0.80}Al_{0.80}Si_{0.20}O_2$ when the Si:Al ratio was increased beyond 1:4. Beyond KAlSiO₄ (which itself may be obtained with a potassium occupancy of less than unity) increase in Si: Al ratio leads to formation of $KAlSi_2O_6$, leucite (13) and finally to KAlSi₃O₈, sanidine. In it, as in all feldspars (13), the alkali ions are contained in large cage sites in the three dimensional framework of aluminosilicate tetrahedra, with one cage site to four tetrahedral sites. At very high pressures, sanidine is transformed to the hollandite form (57), isotypical with KAlTi₃O₈, while silica itself transforms to rutile form (58), which has no sites for additional ions.

Conclusion

The structures examined in the present paper each contain alkali ions in sites contained within a framework of coordination polyhedra formed about ions considerably smaller than sodium or potassium. Such structures appear almost universally to be able to support a nonstoichiometry based on fractional occupancy of the alkali ion sites, with a corresponding charge balance being provided within the fully occupied octahedral or tetrahedral framework. This phenomenon is not of course confined to aluminates, ferrites, or titanates, although the alkali-containing ternary titanates are a particularly rich source of such structures. It is also the basis of the marked nonstoichiometry of the various alkalicontaining "bronzes" of vanadium, molybdenum, and tungsten, as well as the basis of deviations from stoichiometry in the alkali-metal aluminosilicates and of the incorporation of alkali ion impurities in the natural modifications of silica.

Acknowledgments

The authors gratefully acknowledge computer programs supplied by Mr. C. H. J. Johnson, Division of Applied Chemistry, C.S.I.R.O., and by Dr. J. J. Daly and Dr. P. J. Wheatley, Monsanto S.A., Zurich (59).

References

 R. W. G. WYCKOFF, "Crystal Structures," 2nd ed., Vol. 2. Wiley, New York, 1964.

- R. HOPPE, "Mem. du Coll. Int. Cent. Nat. Rech. Sci.," Bull. Soc. Chim. Fr., 5th ser, 1115–1121 (1965).
- V. I. SPITSYN, I. A. MURAV'EVA, L. M. KOVBA, AND I. I. KORCHAK, Russ. J. Inorg. Chem. 14, 759 (1969).
- 4. S. GOLDSZTAUB, C. R. H. Acad. Sci. 196, 280 (1933).
- 5. S. HILPERT AND A. LINDNER, Z. Phys. Chem. Abt. B 22, 395 (1933).
- 6. F. BERTAUT AND P. BLUM, C. R. H. Acad. Sci. 239, 429 (1954).
- 7. J. Théry, D. BRIANCON, AND R. COLLONGUES, C. R. H. Acad. Sci. 252, 1475 (1961).
- 8. E. F. BERTAUT, A. DELAPALME, AND G. BASSI, C. R. H. Acad. Sci. 257, 421 (1963).
- A. F. REID, A. D. WADSLEY, AND M. J. SIENKO, *Inorg. Chem.* 7, 112 (1968).
- 10. W. G. MUMME AND A. F. REID, Acta Crystallogr. Sect. B 24, 625 (1968).
- 11. G. F. CLARINGBULL AND F. A. BANNISTER, Acta Crystallogr. 1, 42 (1948).
- 12. A. J. PEROTTA AND J. V. SMITH, Acta Crystallogr. 16, A13 (1963).
- 12. A. J. PEROTTA AND J. V. SMITH, Acta Crystallogr. Sect. A 16, A13 (1963).
- 13. L. BRAGG AND G. F. CLARINGBULL, "Crystal Structures of Minerals," Bell, London, 1965.
- 14. W. RÜDORFF AND H. BECKER, Z. Naturforsch. 96, 614 (1954).
- E. KORDES AND E. RÖTTIG, Z. Anorg. Allg. Chem. 263, 34 (1951).
- 16. A. F. REID AND A. E. RINGWOOD, Inorg. Chem. 7, 443 (1968).
- E. KORDES AND J. PETZOLDT, Z. Anorg. Allg. Chem. 335, 138 (1965).
- 18. R. CLOS, M. DEVALETTE, C. FOUASSIER AND P. HAGENMULLER, Mater. Res. Bull. 5, 179 (1970).
- 19. R. D. SHANNON AND C. T. PREWITT, Acta Crystallogr. Sect. B 25, 925 (1969).
- 20. W. G. MUMME, A. F. REID, AND C. LI, to be published.
- P. HAGENMULLER, A. LECERF, AND M. ONILLON, C. R. H. Acad. Sci. 255, 928 (1962).
- 22. M. DANOT, A. LEBLANC AND J. ROUXEL, Bull. Soc. Chim. Fr., 2670 (1969).
- 23. G. P. ESPINOSA, Inorg. Chem. 3, 848 (1964).
- 24. J. ITO AND C. FRONDEL, Amer. Mineral. 52, 773 (1967).
- S. GELLER, Z. Kristallogr., Kristallgeometrie, Kristallphys., Kristallchem. 125, 1 (1967).
- P. M. HILL, H. S. PEISER AND J. R. RAIT, Acta Crystallogr. 9, 981 (1956).
- E. F. BERTAUT, P. BLUM, AND G. MAGNANO, Bull. Soc. Fr. Mineral. Cristallogr. 129, 536 (1956).
- B. F. DECKER AND J. S. KASPER, Acta Crystallogr. 10, 332 (1957).
- A. D. WADSLEY, Z. Kristallogr. Kristallgeometrie, Kristallphys. Kristallchem. 120, 398 (1964).
- G. BAYER AND W. HOFFMANN, Z. Kristallogr. Kristallgeometrie, Kristallophys., Kristallchem. 121, 1 (1965).
- 31. A. F. REID AND M. J. SIENKO, Inorg. Chem. 6, 321 (1967).
- 32. S. ANDERSSON AND A. D. WADSLEY, Acta Crystallogr. 15, 201 (1962).

- 33. G. TEUFER AND A. K. TEMPLE, Nature (London) 211, 179 (1966).
- 34. I. E. GREY, W. G. MUMME, AND A. F. REID, in press.
- 35. C. ROMERS, C. J. M. ROOYMANS, AND R. A. G. DE GRAAF, Acta Crystallogr. 22, 760 (1967).
- 36. S. ANDERSSON AND A. D. WADSLEY, Acta Crystallogr. 14, 1245 (1961).
- 37. S. ANDERSSON AND A. D. WADSLEY, Acta Crystallogr. 15, 194 (1962).
- 38. A. M. LEVIN, C. R. ROBBINS, AND H. F. MCMURDIE, "Phase Diagrams for Ceramists," 1969 Supplement, American Ceramic Society, 1969.
- 39. G. LANG, Z. Anorg. Allg. Chem. 276, 77 (1954).
- I. N. BELYAEV AND A. G. BELYAEVA, Russ. J. Inorg. Chem. 10, 252 (1965).
- 41. E. K. BELYAEV, N. M. PANASENKO, AND E. V. LINNIK, Russ. J. Inorg. Chem. 15, 336 (1970).
- 42. M. HAMELIN, Bull. Soc. Chim. Fr., 1421 (1957).
- L. PAULING, Z. Kristallogr. Kristallgeometrie, Kristallphys. Kristallchem. 75, 88 (1930).
- 44. T. F. W. BARTH, J. Chem. Phys. 3, 323 (1933).
- R. W. G. WYKOFF, "Crystal Structures," 2nd ed., Vol. 1, Wiley, New York, 1963.
- R. B. SOSMAN, "The Phases of Silica," Rutgers Univ. Press, New Brunswick, N.J., 1965.

- P. B. MOORE AND S. J. LOUISNATHAN, Z. Kristallogr. Kristallgeometrie, Kristallphys. Kristallchem. 130, 438 (1969).
- 48. S. ANDERSSON AND A. D. WADSLEY, Acta Chem. Scand. 15, 663 (1961).
- 49. G. E. BROWN AND G. V. GIBBS, Amer. Mineral. 54, 1528 (1969).
- 50. J. B. JONES, Acta Crystallogr. Sect. B 24, 355 (1968).
- 51. R. D. SHANNON AND C. T. PREWITT, J. Inorg. Nucl. Chem. 32, 1427 (1970).
- 52. H. G. HUCKENHOLZ, *Amer. J. Sci.* 267A, 209 (1969); and private communication by H. G. Huckenholz.
- 53. J. A. WATTS, A. F. REID, AND W. G. MUMME, in press.
- 54. A. F. REID, W. G. MUMME, AND A. D. WADSLEY, Acta Crystallogr. Sect. B 24, 1228 (1968).
- 55. A. BYSTROM AND A. M. BYSTROM, Acta Crystallogr. 3, 146 (1950).
- 56. G. BAYER AND W. HOFFMANN, Amer. Min. 51, 511 (1966).
- 57. A. E. RINGWOOD, A. F. REID, AND A. D. WADSLEY, Acta Crystallogr. 23, 1093 (1967).
- 58. S. M. STISHOV AND N. V. BELOV, Dokl. Akad. Nauk. SSR 143, 951 (1962).
- 59. J. J. DALY, F. S. STEPHENS, AND P. J. WHEATLEY, MONSANTO Research S.A., Final Report No. 52, 1963.